

## Two-micron lasing in $\text{NaLa}_{1/2}\text{Gd}_{1/2}(\text{WO}_4)_2$ crystals doped with $\text{Tm}^{3+}$ ions

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**Abstract.** Lasing on the  ${}^3\text{F}_4 \rightarrow {}^3\text{H}_6$  transition of  $\text{Tm}^{3+}$  ions in  $\text{Tm}^{3+}:\text{NaLa}_{1/2}\text{Gd}_{1/2}(\text{WO}_4)_2$  crystals pumped by a diode laser is obtained for the first time. The  $\pi$ - and  $\sigma$ -polarised laser radiation at wavelengths of 1908 and 1918 nm was generated with a slope efficiency of 28% and 25%, respectively.

**Keywords:** two-micron lasing,  $\text{Tm}^{3+}$  ions,  $\text{NaLa}_{0.5}\text{Gd}_{0.5}(\text{WO}_4)_2$  crystals.

The use of two-micron laser radiation in medicine, as well as in monitoring of various gases (for example,  $\text{NO}_2$ ,  $\text{CO}_2$ ,  $\text{NH}_3$ ) in the atmosphere, stimulates the search for new active media emitting at this wavelength. Two-micron lasing can be obtained in solid matrices doped with  $\text{Ho}^{3+}$  ( ${}^5\text{I}_7 \rightarrow {}^5\text{I}_8$  transition) or  $\text{Tm}^{3+}$  ( ${}^3\text{F}_4 \rightarrow {}^3\text{H}_6$  transition). According to the energy level diagram of  $\text{Tm}^{3+}$  ions in crystalline matrices, the upper laser level  ${}^3\text{F}_4$  of this ion is populated due to the cross-relaxation ( ${}^3\text{H}_4 \rightarrow {}^3\text{F}_4$ ,  ${}^3\text{H}_6 \rightarrow {}^3\text{F}_4$ ) under pumping of the  ${}^3\text{H}_4$  level. This scheme of population of the  ${}^3\text{F}_4$  level is used to obtain two-micron lasing under both lamp pumping [1–3] and pumping by commercial laser diodes ( $\lambda_p \approx 800$  nm) [4, 5].

At present, in order to create compact mid-power two-micron lasers, one studies the crystals of scheelite-like double tungstates doped with  $\text{Tm}^{3+}$  ions [6–12]. The absorption and fluorescence lines of  $\text{Tm}^{3+}$  in these laser hosts are strongly inhomogeneously broadened due to the disordered crystal structure. The broad absorption bands of the active ions make them low-sensitive to variations in the pump source spectrum, while the broad luminescence bands allow one to obtain frequency-tunable laser radiation, including fine tuning to the absorption lines of particular gases.

Lasing in  $\text{Tm}^{3+}:\text{NaLaGd}(\text{WO}_4)_2$  crystals was reported for the first time in [8], and somewhat later the authors of paper [6] achieved tunable lasing in this crystal in the region of  $\sim 200$  nm under pumping by a  $\text{Ti}^{3+}:\text{Al}_2\text{O}_3$  laser, as well as lasing at 1997 nm under diode laser pumping. The authors of [7] reported two-micron lasing in  $\text{Tm}^{3+}:\text{NaLa}(\text{WO}_4)_2$  crystals pumped by a  $\text{Ti}^{3+}:\text{Al}_2\text{O}_3$  laser.

In this paper, we present the results of the first, to our knowledge, experiment on lasing in  $\text{Tm}^{3+}:\text{NaLa}_x\text{Gd}_{1-x}(\text{WO}_4)_2$  crystals. The crystals of mixed scheelite-like tungstates  $\text{Tm}^{3+}:\text{NaLa}_x\text{Gd}_{1-x}(\text{WO}_4)_2$  with a variable La–Gd composition occupy an intermediate position between the lanthanum and gadolinium tungstates. Our estimates show that the absorption cross section for the  ${}^3\text{H}_6 \rightarrow {}^3\text{H}_4$  transition and the stimulated emission cross section for the  ${}^3\text{F}_4 \rightarrow {}^3\text{H}_6$  transition of  $\text{Tm}^{3+}$  ions in mixed scheelite crystals are comparable with the corresponding characteristics for  $\text{Tm}^{3+}:\text{NaGd}(\text{WO}_4)_2$  [6, 9, 12] and  $\text{Tm}^{3+}:\text{NaLa}(\text{WO}_4)_2$  [7] crystals.

The  $\text{Tm}^{3+}:\text{NaLa}_{1/2}\text{Gd}_{1/2}(\text{WO}_4)_2$  crystals were grown by the Czochralski method from a melt of the composition  $\text{NaLa}_{0.46}\text{Gd}_{0.46}\text{Tm}_{0.07}(\text{WO}_4)_2$ . To optimise the pump radiation absorption, the cylindrical active elements (AEs) were cut from a crystal oriented with a DRON-4 X-ray diffractometer so that the optical axis of the crystal was perpendicular to the axis of the cylindrical AE. The AE dimensions (diameter 3 mm, length 5 mm) satisfied the condition of efficient longitudinal diode laser pumping at  $\lambda_p \sim 794$  nm.

The AE faces were AR-coated for the pump and laser wavelengths. The residual reflectance at the pump ( $R_{794}$ ) and laser ( $R_{1910}$ ) wavelengths was 0.67% and 0.21%, respectively.

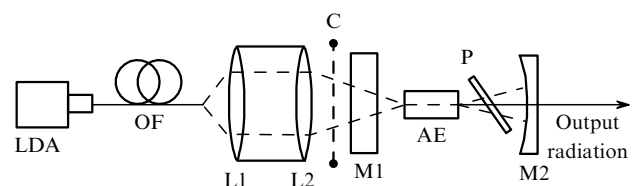
The scheme used to obtain lasing in the  $\text{Tm}^{3+}:\text{NaLa}_{1/2}\text{Gd}_{1/2}(\text{WO}_4)_2$  crystal is shown in Fig. 1.

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**Figure 1.** Laser scheme: (LDA) laser diode array, (OF) optical fibre, (L1) and (L2) telescope lenses, (C) chopper, (M1) dichroic mirror, (AE) active element, (P) glass plate, (M2) plano-concave mirror.

As a pump source, we used a 40-W laser diode array (LDA), whose temperature was kept constant at  $\sim 26.5^\circ\text{C}$ , which corresponded to the radiation wavelength of  $\sim 794\text{ nm}$  (the diode radiation linewidth was  $\sim 2\text{ nm}$ ). To decrease the heat load on the AE, the average pump power was attenuated using a chopper, which formed radiation pulses with a duration of 10 ms and a repetition rate of 5 Hz. The radiation of the fibre-coupled LDA (fibre diameter 800  $\mu\text{m}$ , numerical aperture 0.14) was focused into the AE by a two-lens telescope. For efficient cooling, the AE wrapped with indium foil was mounted in a copper heatsink, whose temperature was kept constant at  $\sim 18^\circ\text{C}$  with the help of a Peltier element. The cavity used in the experiment was formed by a plane dichroic mirror M1 ( $T_p \approx 90\%$ ,  $T_{\text{las}} \approx 0.5\%$ ) and a plano-concave mirror M2 ( $T_{\text{las}} \approx 11\%$ ) with a 200-mm radius of curvature of the spherical face. To separate laser radiation with the  $\pi$  or  $\sigma$  polarisation, we placed a glass plate inside the cavity at Brewster angle to the system axis.

During the experiment with the  $\text{Tm}^{3+}:\text{NaLa}_{1/2}\text{Gd}_{1/2}(\text{WO}_4)_2$  crystal, we obtained laser radiation with the  $\pi$  and  $\sigma$  polarisations according to the spectral dependences of the stimulated emission cross section for the  ${}^3\text{F}_4 \rightarrow {}^3\text{H}_6$  transition of  $\text{Tm}^{3+}$  in the  $\text{Tm}^{3+}:\text{NaGd}(\text{WO}_4)_2$  [6],  $\text{Tm}^{3+}:\text{NaLa}(\text{WO}_4)_2$  [7], and  $\text{Tm}^{3+}:\text{NaLa}_{1/2}\text{Gd}_{1/2}(\text{WO}_4)_2$  crystals for the  $\pi$  and  $\sigma$ -polarisations. The lasing characteristics of the  $\text{Tm}^{3+}:\text{NaLa}_{1/2}\text{Gd}_{1/2}(\text{WO}_4)_2$  crystal for these polarisations are given in Table 1.

**Table 1.** Lasing parameters of the  $\text{Tm}^{3+}:\text{NaLa}_{1/2}\text{Gd}_{1/2}(\text{WO}_4)_2$  crystal.

Polarisation	Slope efficiency (%)	Total efficiency (%)	Laser wavelength/nm
$\pi$	28	11	1908
$\sigma$	25	9	1918

Figure 2 shows the dependences of the output laser power on the pump power for the  $\pi$  and  $\sigma$  polarisations. The slope efficiency was 25 % for the  $\sigma$  polarisation and 28 % for the  $\pi$  polarisation.

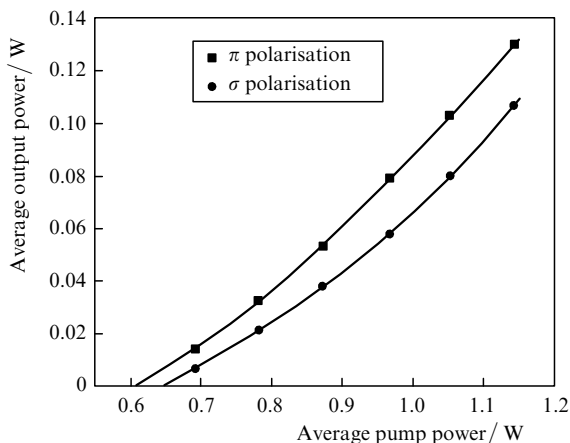
Our estimates of the spectroscopic characteristics of  $\text{Tm}^{3+}$  ions in scheelite crystals with a varying La–Gd composition together with the results of the reported experiment allow us to conclude that, using AEs of a better

optical quality and optimising the output mirror transmittance and the cavity scheme, it is possible to achieve higher output laser characteristics for the  $\text{Tm}^{3+}:\text{NaLa}_{1/2}\text{Gd}_{1/2}(\text{WO}_4)_2$  crystal. In addition, the broad gain bands in the region of 2  $\mu\text{m}$  for the  $\pi$ - and  $\sigma$ -polarised radiation in this crystal allow one to obtain wavelength-tunable laser radiation.

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**Figure 2.** Dependences of the average laser output power on the average pump power.