LASERS

PACS numbers: 42.55.Ks DOI: 10.1070/QE2014v044n12ABEH015603

Mechanism of singlet oxygen deactivation in an electric discharge oxygen-iodine laser

V.N. Azyazov, P.A. Mikheyev, A.A. Pershin, A.P. Torbin, M.C. Heaven

Abstract. We have determined the influence of the reaction of molecular singlet oxygen with a vibrationally excited ozone molecule $O_2(a^{1}\Delta) + O_3(v) \rightarrow 2O_2 + O$ on the removal rate of $O_2(a^{1}\Delta)$ in an electric-discharge-driven oxygen-iodine laser. This reaction has been shown to be a major channel of $O_2(a^{1}\Delta)$ loss at the output of an electric-discharge singlet oxygen generator. In addition, it can also contribute significantly to the loss of $O_2(a^{1}\Delta)$ in the discharge region of the generator.

Keywords: singlet oxygen, oxygen–iodine laser, deactivation, recombination, vibrationally excited ozone, discharge, $O_2(a^{l}\Delta)$, O_3 , $O({}^{3}P)$.

Small-signal gain of the active medium of an electric-discharge-driven oxygen–iodine laser (EOIL) [1] does not make it possible to effectively extract the energy stored in singlet oxygen $O_2(a^{1}\Delta)$ [2]. To increase the gain, one should increase the concentration of both iodine atoms and $O_2(a^{1}\Delta)$ molecules. Vasiljeva et al. found [3] that the rate of $O_2(a^{1}\Delta)$ deactivation at the output of an electric-discharge singlet oxygen generator (ED SOG) increases with increasing concentrations of oxygen atoms [O], molecular oxygen [O₂] and the buffer gas [M]. The rate of $O_2(a^{1}\Delta)$ loss at the ED SOG output is satisfactorily explained by the phenomenological three-body deactivation mechanism [3]

$$O + O_2(a^1\Delta) + M \rightarrow O + O_2 + M.$$
⁽¹⁾

Azyazov et al. [4-6] also observed an abnormally high rate of O₂(a¹ Δ) deactivation after laser photolysis of O-O₂-Ar-He-CO₂ mixtures. In this case, the addition of Ar had no effect on the rate of O₂(a¹ Δ) deactivation, whereas the addition of He and CO₂ even reduced it [6], which is in contradiction with mechanism (1). The rate of O₂(a¹ Δ) loss

Received 27 June 2014; revision received 10 September 2014 *Kvantovaya Elektronika* **443** (12) 1083–1084 (2014) Translated by I.A. Ulitkin after laser photolysis is well explained by the chemical process [5, 6]

$$O_2(a^1\Delta) + O_3(v) \rightarrow 2O_2 + O, \tag{2}$$

where vibrationally excited ozone $O_3(v)$ is produced during the three-body recombination [7]

$$O + O_2 + M \rightarrow O_3(v) + M.$$
(3)

Here, $v = v_1 + v_2 + v_3$ is the total number of quanta at deformation ($v_2 = 701 \text{ cm}^{-1}$), symmetrical ($v_1 = 1103 \text{ cm}^{-1}$) and antisymmetrical ($v_3 = 1042 \text{ cm}^{-1}$) stretching modes of ozone [7]. The authors of papers [7–10] also found that $O_3(v)$ efficiently reacts with $O_2(a^{1}\Delta)$. Despite this, process (2) has not previously been used to explain the high rate of $O_2(a^{1}\Delta)$ loss at the ED SOG output. The aim of this study is to determine the effect of process (2) on the dynamics of $O_2(a^{1}\Delta)$ in an EOIL.

The kinetic scheme of the processes involving vibrationally excited ozone is given in [6]. The absence of the measured probabilities of excitation of vibrational modes of the O₃ molecule in process (3) and of the rate constants of process (2) for specific sets of v_1 , v_2 and v_3 values complicates the modelling of vibrational kinetics of ozone. Azyazov and Heaven [5] showed that the simplified model of vibrational kinetics of ozone with a combined mode v adequately describes the experimental results obtained in Refs [4–6]. In process (3) an ozone molecule with $v \ge 3$ is produced [5, 7]. In collisions with the particles of the medium, vibrational quanta are redistributed among the three modes. The ozone molecule O₃(v) is deactivated in the VT-process [7]

$$O_3(v) + M \to O_3(v-1) + M$$
 (4)

or removed in chemical reactions (2) and

$$O_3(v) + O \rightarrow O_2 + O_2. \tag{5}$$

Reactions (2) and (5) with v < 2 are slow and have no significant effect on the kinetics of O₃(v) [5]. For $v \ge 2$ Azyazov and Heaven [5] give the following values of the rate constants of reactions (2) and (5): $k_2 = 4.1 \times 10^{-11}$ cm³ s⁻¹ and $k_5 = 1.2 \times 10^{-11}$ cm³ s⁻¹.

Quasi-stationary concentration of $O_3(v \ge 2)$ can be obtained from the balance of its formation in process (3) and loss in processes (2), (4) and (5):

$$[O_3(v \ge 2)] = \frac{\sum_{M} k_3^M[O][O_2][M]}{k_2[O_2(a)] + \sum_{M} k_4^M[M] + k_5[O]}.$$
 (6)

^{V.N. Azyazov, A.P. Torbin Samara Branch of P.N. Lebedev Physics} Institute, Russian Academy of Sciences, ul. Novo-Sadovaya 221, 443011 Samara, Russia; Samara State Aerospace University, Moskovskoe shosse 34, 443086 Samara, Russia; e-mail: azyazov@fian.smr.ru, torbinalex@gmail.com;
P.A. Mikheyev Samara Branch of P.N. Lebedev Physics Institute, Russian Academy of Sciences, ul. Novo-Sadovaya 221, 443011 Samara, Russia; e-mail: mikheyev@fian.smr.ru;
A.A. Pershin Samara State Aerospace University, Moskovskoe shosse 34, 443086 Samara, Russia; e-mail: anchizh93@gmail.com;
M.C. Heaven Emory University, Atlanta, GA, 30322, USA; e-mail: mheaven@emory.edu

Let us show that process (2) provides the same rates of $O_2(a^1\Delta)$ loss as process (1) under the experimental conditions of paper [3], where the process of three-body deactivation was first proposed. To this end, we consider the ratio of the rates of these processes, taking into account (6):

$$Q_{2/1} = \frac{k_2[O_3(v \ge 2)][O_2(a)]}{\sum_{\mathbf{M}} k_1^{\mathbf{M}}[O][O_2(a)][\mathbf{M}]}$$
$$= \frac{k_2[O_2]\sum_{\mathbf{M}} k_3^{\mathbf{M}}[\mathbf{M}]}{\sum_{\mathbf{M}} k_1^{\mathbf{M}}[\mathbf{M}] (k_2[O_2(a)] + \sum_{\mathbf{M}} k_4^{\mathbf{M}}[\mathbf{M}] + k_5[O])}.$$

Under the experimental conditions of paper [3] (Ar:O₂ = 99:1; mixture pressure, 100 Torr; concentration of singlet oxygen molecules, $[O_2(a)] = 1.5 \times 10^{15} \text{ cm}^{-3}$ and $[O] = 2 \times 10^{15} \text{ cm}^{-3}$; and mixture temperature, T = 300 K) at $k_1^{\text{Ar}} = 0.62 \times 10^{-32} \text{ cm}^6 \text{ s}^{-1}$ [3], $k_1^{\text{O2}} = 1 \times 10^{-32} \text{ cm}^6 \text{ s}^{-1}$ [3], $k_3^{\text{O2}} = 6 \times 10^{-34} \text{ cm}^6 \text{ s}^{-1}$ [3], $k_3^{\text{Ar}} = 0.62 \times k_3^{\text{O2}}$ [3], $k_4^{\text{O2}} = 3 \times 10^{-14} \text{ cm}^3 \text{ s}^{-1}$ [7] and $k_4^{\text{Ar}} = 5.9 \times 10^{-15} \text{ cm}^3 \text{ s}^{-1}$ [7], the ratio $Q_{2/1}$ is close to unity, i.e., the rate of $O_2(a^{1}\Delta)$ loss at the ED SOG output is equally provided by processes (1) and (2).

The rate of process (2) with (6) taken into account can be represented as:

$$R_{2} = k_{2}[O_{3}(v \ge 2)][O_{2}(a)] = \frac{k_{2}[O_{2}(a)]\sum_{M} k_{3}^{M}[O][O_{2}][M]}{k_{2}[O_{2}(a)] + \sum_{M} k_{4}^{M}[M] + k_{5}[O]}$$

For typical conditions in a post-discharge zone of the ED SOG [1] ([He] = 9×10^{17} cm⁻³, [O₂] = 2.5×10^{17} cm⁻³, [O] = 5×10^{15} cm⁻³, [O₂(a)] = 3×10^{16} cm⁻³ and T = 550 K) the terms in the denominator satisfy the condition

$$k_2[O_2(\mathbf{a})] \gg \sum_{\mathbf{M}} k_4^{\mathbf{M}}[\mathbf{M}] + k_5[\mathbf{O}].$$

In this case,

$$R_2 \approx \sum_{\mathbf{M}} k_3^{\mathbf{M}}[\mathbf{O}][\mathbf{O}_2][\mathbf{M}],$$

and the rate of process (2) is limited by the rate of $O_3(v)$ production during three-body recombination (3). This was the reason why the authors of [3] considered the three-body process to be main mechanism of $O_2(a^{1}\Delta)$ deactivation.

Reaction (2) can also make a significant contribution to the rate of $O_2(a^1\Delta)$ loss in the discharge zone of the ED SOG. Consider the ratio of the rate of this reaction to the rate of the fastest process in the discharge zone – deactivation of $O_2(a^1\Delta)$ by an electron impact:

$$O_2(a^1\Delta) + e \to O_2 + e. \tag{7}$$

Under the experimental conditions of paper [11], the ratio of the rates of processes (2) and (7), $Q_{2/7}$, is ~0.2 in the case when O atoms are not removed from the system. In experiments with a reduced concentration of O atoms, the ratio is much smaller, $Q_{2/7} = 0.03$, due to addition of NO and coating of the chamber walls with mercury oxide. Consequently, process (2) also makes a significant contribution to the rate of $O_2(a^{1}\Delta)$ loss in the discharge zone of the ED SOG in the case of excess O atoms. Braginsky et al. [11] produced a record concentration of $O_2(a^{1}\Delta)$ in the ED SOG by removing excess oxygen atoms. Thus, process (2) can ensure the observed rates of $O_2(a^{1}\Delta)$ deactivation both in the post-discharge zone [3] and after laser photolysis [5]. Processes (2)–(5) should also be included in the kinetic scheme of the processes in the discharge zone of the ED SOG. The authors of [4–9] present experimental evidence to support a decisive role of process (2) in deactivation of singlet oxygen, whereas process (1) is phenomenological and cannot provide the rate of $O_2(a^{1}\Delta)$ loss after laser photolysis [4, 5]. The rate of $O_2(a^{1}\Delta)$ deactivation can be reduced by removing excess O atoms, for example, by adding NO into the mixture [1] or by coating the chamber walls with mercury oxide [11] and by adding $O_3(v)$ quenchers, such as CO₂, SF₆, SiF₄, etc. to the mixture at the ED SOG output.

Acknowledgements. This work at the Samara State Aerospace University (SSAU) was supported by the Ministry of Education of the Russian Federation in the framework of the programme aimed at increasing competitiveness of SSAU in 2013–2020 and by the State Task No. 3.161.2014/K for Universities and Research Organisations in the field of scientific activity; the work in the Samara branch of Lebedev Physics Institute was supported by the Russian Foundation for Basic Research (Grant No. 14-05-97013).

References

- Benavides G.F., Woodard B.S., Zimmerman J.W., Palla A.D., Day M.T., King D.M., Carroll D.L., Verdeyen J.T., Solomon W.C. *IEEE J. Quantum Electron.*, 48, 741 (2012).
- Mezhenin A.V., Azyazov V.N. IEEE J. Quantum Electron., 49, 739 (2013).
- Vasiljeva A.N., Klopovskiy K.S., Kovalev A.S., Lopaev D.V., Mankelevich Y.A., Popov N.A., Rakhimov A.T., Rakhimova T.V. J. Phys. D: Appl. Phys., 37, 2455 (2004).
- Azyazov V.N., Mikheyev P.A., Postell D., Heaven M.C. Chem. Phys. Lett., 482, 56 (2009).
- Azyazov V.N., Heaven M.C. Int. J. Chem. Kinet. (2014, in press).
 Azyazov V.N., Mikheyev P.A., Heaven M.C. Proc. SPIE Int.
- Azyazov V.N., Mikheyev P.A., Heaven M.C. Proc. SPIE Int. Soc. Opt. Eng., 7751, 77510E (2010).
- Steinfield J.I., Adler-Golden S.M., Gallagher J.W. J. Phys. Chem. Ref. Data, 16, 911 (1987).
- Kurylo M.J., Braun W., Kaldor A., Freund S.M., Wayne R.P. J. Photochem., 3, 71 (1974).
- Rawlins W.T., Caledonia G.E., Armstrong R.A. J. Chem. Phys., 87, 5209 (1987).
- Klopovskii K.S., Kovalev A.S., Lopaev D.V., Rakhimov A.T., Rakhimov T.V. *Fiz. Plazmy*, 18, 1606 (1992).
- Braginsky O.V., Kovalev A.S., Lopaev D.V., Proshina O.V., Rakhimova T.V., Rakhimov A.T., Vasilieva A.N. J. Phys. D: Appl. Phys., 40, 6571 (2008).