Tunable 2- μ m lasing in calcium – niobium – gallium garnet crystals doped with Ho³⁺ ions

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Abstract. Lasing on the ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$ transition of Ho³⁺ ions in holmium-doped calcium-niobium-gallium garnet crystals is obtained at a wavelength of about 2095 nm with an output power of 2.1 W under pumping by a laser based on Tm : LiYF₄ crystal. Tunable lasing in these crystals within a wavelength range of 2045-2120 nm is achieved using an interference-polarisation filter.

Keywords: crystals with garnet structure, Ho³⁺ ions, lasing.

1. Introduction

Calcium-niobium-gallium garnet (CNGG) and calciumlithium-niobium-gallium garnet (CLNGG) crystals belong to crystals with disordered crystal structure. These crystals were synthesised for the first time more than thirty years ago [1]. They are characterised by a lower melting temperature (1430-1470°C) than yttrium-aluminium garnet (YAG), gadolinium-scandium-aluminium garnet (GSAG) and gadolinium-scandium-gallium garnet (GSGG) crystals. This simplifies the growth of CNGG and CLNGG crystals and allows one to synthesise them by technologies that do not use iridium.

Due to the disordered crystal structure of CNGG and CLNGG, the absorption and luminescence spectra of rareearth (RE) ions in these crystals are inhomogeneously broadened. This is of interest for developing tunable and modelocked lasers based on CNGG and CLNGG crystals doped with RE ions.

At present, there exist extensive experimental data on the spectral-luminescent and laser properties of CNGG and CLNGG crystals doped with Nd³⁺, Tm³⁺, and Yb³⁺+ ions [2–14]. Some results of investigation of the spectroscopic characteristics of Ho³⁺: CNGG crystals are presented in [15]. However, we failed to find works on the laser properties of these crystals resonantly pumped to the ${}^{5}I_{7}$ level of Ho³⁺ ions in the available literature.

The aim of the present work is to study the laser characteristics of Ho:CNGG crystals under resonant pumping by a Tm: LiYF₄ laser to the ${}^{5}I_{7}$ level of Ho³⁺ ions.

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2. Experimental samples and methods

The calcium–niobium–gallium garnet crystals doped with Ho^{3+} ions [Ho:Ca₃(NbGa)₅O₁₂] were grown by the Czochralski methods in a platinum crucible. The concentration of Ho^{3+} ions in the grown crystals was 0.7 at %. To perform spectral-luminescent measurements, we cut plane-parallel plates 7.5 mm thick from the grown CNGG crystals. In laser experiments, we used a cylindrical active element 18 mm long and 4 mm in diameter.

The absorption spectra were recorded on a Lambda 950 PerkinElmer spectrophotometer. A halogen lamp was used as a radiation source. As a detector for measuring the absorption spectra at the ${}^{5}I_{8} \rightarrow {}^{5}I_{7}$ transition of Ho³⁺ ions, we used a PbS photodetector. The luminescence spectra at this transition were measured using an MDR-23 monochromator. For excitation we used a solid-state MGL-F-532nm-2W laser based on Nd: YVO₄ crystal with a wavelength of 532 nm. The luminescence signal was recorded by a PbS photoresistor.

The laser spectrum was measured using a Horiba HR 1000 spectrometer and a PbS photodetector.

3. Experimental results

To achieve lasing on the ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$ transition under resonant pumping to the ${}^{5}I_{7}$ level of Ho³⁺ ions, one must know the spectral-luminescent characteristics of the ${}^{5}I_{8} \rightarrow {}^{5}I_{7}$ and ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$ optical transitions of Ho³⁺ ions in the corresponding crystal matrix. The absorption spectrum on the ${}^{5}I_{8} \rightarrow {}^{5}I_{7}$ transitions of Ho³⁺ ions in the active element at T = 300 K is shown in Fig. 1.

The wavelength dependence of the stimulated emission cross section $\sigma_{\rm em}(\lambda)$ for the ${}^{5}\mathrm{I}_{7} \rightarrow {}^{5}\mathrm{I}_{8}$ transition of Ho³⁺ ions in a Ho: CNGG crystal ($C_{\rm Ho} = 0.7$ at %) was determined by the Füchtbauer–Ladenburg formula [16]

$$\sigma_{\rm em}(\lambda) = \frac{\lambda^5 \beta_{JJ'} I_{\rm em}(\lambda)}{8\pi c n^2 \tau_{\rm rad} \int I_{\rm em}(\lambda) \lambda d\lambda},\tag{1}$$

where τ_{rad} is the radiative lifetime of the ${}^{5}I_{7}$ level of Ho³⁺ ions; *n* is the refractive index of the medium; $I_{em}(\lambda)$ is the relative luminescence intensity; *c* is the speed of light; $\beta_{JJ'}$ is the luminescence branching ratio; and λ is the wavelength. For the ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$ transition, $\beta_{JJ'} = 1$. The lifetime τ_{rad} was determined as the reciprocal probability of the ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$ transition of Ho³⁺ ions in Ho:CNGG crystals [17] and was found to be 13 ms.

The wavelength dependences of the absorption and stimulated emission cross sections for the ${}^{5}I_{8} \rightarrow {}^{5}I_{7}$ and ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$ transitions of Ho³⁺ ions are presented in Fig. 2.



Figure 1. Absorption spectrum for the ${}^{5}I_{8} \rightarrow {}^{5}I_{7}$ transition of Ho³⁺ ions in the Ho:CNGG active element ($C_{Ho} = 0.7$ at %); T = 300 K. The arrow indicates the wavelength of optical pumping into the ${}^{5}I_{7}$ level of Ho³⁺ ions in Ho:CNGG crystals.



Figure 2. Spectral dependences of the absorption and stimulated emission cross sections on the ${}^{5}I_{8} \rightarrow {}^{5}I_{7}$ and ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$ transitions of Ho³⁺ ions; T = 300 K.

Using the found spectral dependences $\sigma_{abs}(\lambda)$ and $\sigma_{em}(\lambda)$, one can determine the spectral profile of the gain for the ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$ laser transition of Ho³⁺ ions in Ho:CNGG crystals as

$$\sigma_{\rm g}(\lambda) = P \sigma_{\rm em}(\lambda) - (1 - P) \sigma_{\rm abs}(\lambda), \tag{2}$$

where *P* is the relative inverse population of levels. Figure 3 presents the calculated gain profiles σ_g for P = 0.1, 0.2, 0.3. One can see that the gain region at P = 0.3 for Ho:Ca₃(NbGa)₅O₁₂ crystals corresponds to the wavelength range of 1970–2200 nm.

The absorption spectra of Ho³⁺ ions in Ho:CNGG crystals show that the best wavelength for efficient pumping of the active element is $\lambda_p = 1925$ nm (see Fig. 1). Because of this, we used a Tm:LiYF₄ laser with wavelength $\lambda_p = 1925$ nm as a pump source. This wavelength was obtained as a result of intracavity selection with the use of an interference–polarisation filter (IPF). The Tm:LiYF₄ laser was pumped by two fibre-coupled diode arrays (DAs). The maximum output power of the pump laser was 15 W.

The optical scheme of the solid-state Ho:CNGG laser pumped by a Tm:LiYF₄ laser is shown in Fig. 4. The pump beam was focused into a Ho:CNGG active element by lenses L1 and L2 with focal distances of 200 and 100 mm, respectively. The pump beam waist diameter was 160 μ m. The



Figure 3. Spectral profile of the gain for the ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$ transition of Ho³⁺ ions in Ho: CNGG crystals ($C_{Ho} = 0.7$ at %) at inverse populations P = 0.1, 0.2, 0.3.

Ho:CNGG laser cavity was formed by a plane input mirror with transmittance T = 90% at the pump wavelength and reflectance R = 99.8% at the lasing wavelength. As output mirrors, we used spherical mirrors with curvature radius r = 150 mm and transmittances T = 2.1%, 3.9%, and 6.8% at the lasing wavelength. The maximum power and efficiency of the Ho:CNGG laser was achieved using an output mirror with transmittance T = 6.8%.



Figure 4. Optical scheme of a solid-state Ho:CNGG laser pumped by a $Tm:LiYF_4$ laser.

Figure 5 presents the Ho: CNGG laser spectrum. One can see that the laser operates at a wavelength of about 2095 nm.

Figure 6 shows the dependence of the output Ho:CNGG laser power on the pump power absorbed in the active element in the case of the output mirror with transmittance T = 6.8%. In this case, the active element absorbed less than 50% of the incident pump power. The lasing threshold with respect to the absorbed power was $P_{\rm th} = 0.7$ W. The slope efficiency with respect to the pump power absorbed in the active element was $\eta = 37\%$. It should be noted that this efficiency was achieved for the active element without antireflection coatings on faces.

Using an interference – polarisation filter (an Al_2O_3 planeparallel plate 4 mm thick), we tuned the wavelength of the



Figure 5. Spectrum of the Ho: CNGG laser at a laser power of 1.5 W.



Figure 6. Dependence of the Ho:CNGG laser power on the pump power absorbed in the active element.



Figure 7. Tuning curve of the Ho: CNGG laser at an incident pump power of 8.1 W.

Ho:CNGG laser. It is found that the spectral range of the Ho:CNGG laser operation is 2045–2120 nm (Fig. 7).

4. Conclusions

We achieved lasing on the ${}^{5}I_{7} \rightarrow {}^{5}I_{8}$ transition of Ho³⁺ ions in crystals of calcium-niobium-gallium garnet doped with Ho³⁺ ions under pumping by a Tm:LiYF₄ laser ($\lambda_{p} = 1925$ nm). The Ho:CNGG operated at a wavelength of about 2095 nm with an output power of 2.1 W.

Placing an interference–polarisation filter in the Ho:CNGG laser cavity, we obtained tunable lasing within the wavelength range of 2045–2120 nm, which testifies to the possibility of fabricating an efficient laser emitting in this wavelength range.

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